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The Crystal Structure of Copper Glutamate Dihydrate*

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The crystal structure of copper(II) glutamate dihydrate, $CuC_5H_7NO_4.2H_2O$, has been determined and refined by three-dimensional least-squares methods. The crystals are orthorhombic, space group $P2_12_12_1$, with a=11.084, b=10.350, c=7.238 Å and four molecules per unit cell. Intensity data were collected visually from Weissenberg photographs about all three crystal axes; of 1087 reflections within the effective sphere of copper radiation all but 16 were strong enough to be measured. The final R index is 0.032; the standard deviations are about 0.004 Å for the light-atom positions.

The coordination about the copper atom is approximately square planar, involving two oxygen atoms and a nitrogen atom of glutamate groups and a water molecule; the Cu-O and Cu-N distances range from 1.97 to 2.00 Å. Two additional glutamate oxygen atoms at 2.30 and 2.59 Å complete a severely distorted octahedron. Five of the six available protons are involved in hydrogen bonding. The glutamate group is in an extended configuration and has normal bond distances and angles.

Introduction

The recent surge of interest in complexes of the transition metal elements with amino acids and peptides has prompted this investigation of the crystal structure of copper(II) glutamate dihydrate, CuC₅H₇NO₄.2H₂O.

Experimental

Crystals were readily grown by evaporation of an aqueous solution of glutamic acid and cupric nitrate. They are blue-green prisms, elongated along c, with principal faces {110} and {011}. A preliminary set of Weissenberg photographs indicated them to be orthorhombic with space group $P2_12_12_1$ (systematic absences: h00, 0k0, and 00l with h, k or l odd). The unitcell dimensions were obtained from a least-squares treatment of measurements on zero-level Weissenberg photographs taken about the three crystallographic axes; the film was held in the asymmetric position. A different crystal was used for each photograph. Absorption and eccentricity coefficients were included as parameters in the least-squares calculations, and weights were assigned inversely proportional to $\sin^2 2\theta$ (reflecting the expected pattern of uncertainties in the observational quantities $\sin^2\theta$). The resulting cell dimensions and their standard deviations are given in Table 1.

The density measured by flotation is 1.954 g.cm⁻³; that calculated on the basis of four molecules of the dihydrate per unit cell is 1.957 g.cm⁻³.

The differences between the cell dimensions for the three crystals (Table 1) are far greater than the experi-

Table 1. Unit-cell dimensions

The cell dimensions and their standard deviations as determined by least-squares treatment of zero-level Weissenberg data for three different crystals are given in the first three columns; the values selected as the best average dimensions are given in the last column.

	No. 1	No.2	No.3	'Best'
a	11.0899 (1)		11.0781 (1)	11.084 (6)
b	,	10.3523 (4)	10.3468 (2)	10.350 (3)
c	7.2398 (1)	7.2358 (2)		7.238 (2)
	(Cu F	Kα ₁ , 1·54051; Cι	$K\alpha_2$, 1.54433)	

mental uncertainties, and are almost surely real. We have no explanation as to the source of the differences. The temperature of the room where the photographs were prepared was approximately constant, a check made during the preparation of the hk0 and 0kl photographs showing no variation exceeding 1° from the average temperature of 24.5°C. It is possible that the amount of water of crystallization varied slightly from one crystal to another; however, no deficiency in the occupancies of the sites of the water molecules was evidenced either by the density measurement or by the refinement of the structural parameters.

For the determination and preliminary refinement of the structure, intensity data were collected from multiple-film equi-inclination Weissenberg photographs of layers 0-6 about c, using a needle-shaped crystal about 0·1 mm in diameter, and of the zero layer about b, using a rhomb-shaped crystal also about 0·1 mm in diameter. Copper radiation was used, and the intensities were estimated visually and corrected for Lorentz and polarization factors. No absorption or extinction corrections were made.

Prior to the final refinement of the structure, additional data were collected in a similar manner for layers 0-9 about the a and b axes. Data about all three axes were then processed through a subprogram of the CRYRM system on an IBM 7094 computer (Duchamp, 1964). This processing included the assignment of esti-

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mated standard deviations to the individual observations, evaluation of film and pack factors relating different photographs of the same layer line, correction for Lorentz and polarization effects, evaluation of relative scale factors for the 27 different sets of photographs, and, finally, the scaling together of these different photographs to give a single value of F^2 , and its standard deviation, for each of the 1087 reflections within the effective sphere of copper radiation. Of these 1087 reflections, only 16 were too weak to be observed.

Determination and refinement of the structure

The trial structure was derived by the usual Patterson and Fourier methods based on the hk0 and h0l data; the Cu-Cu vectors were pronounced features of the Patterson projection onto (001). Preliminary refinement was by difference projections; the first three-dimensional structure-factor calculation led to an R index of 0.10.

The initial three-dimensional refinement was based on the data collected about the c axis only. Block-diagonal least-squares refinement, carried out on a Burroughs 220 computer using a program written by A. Hybl, reduced R to 0.052; anisotropic temperature factors were assigned to the heavy atoms, and the five hydrogen atoms of the carbon chain – whose positions could be assigned from structural considerations – were included in the structure factors but were not refined. The remaining hydrogen atoms were eventually located with the help of difference maps, and further refinement of the heavy-atom parameters led to an R index of 0.040. The results at this stage were presented at the International Congress in Rome (Gramaccioli, 1963).

Final refinement was carried out on an IBM 7094 computer, and was based on the complete set of intensities measured about all three axes (see *Experimental*). Refinement was by full-matrix least squares, using the orthorhombic subprogram of the CRYRM system (Duchamp, 1964). The quantity minimized was Σw $(F_o^2 - 1/k^2 F_c^2)^2$, with weights w taken equal to the reciprocals of the squares of the standard deviations in the individual values of F_o^2 as derived during the data-reduction process. For 38 of the strongest reflections presumed to suffer from extinction and 9 very weak reflections whose intensity estimations were doubtful, the weight was taken to be 0. Atomic form factors were taken from International Tables for X-ray Crystallography (1962), the values for Cu being corrected for the real part $(\Delta F')$ of the anomalous dispersion effect. A total of 162 parameters were adjusted: coordinates of all 24 atoms in the molecule, anisotropic temperature parameters of the 13 heavier atoms, isotropic temperature factors of the hydrogen atoms, and a scale fac-

Considerable difficulty was encountered in locating the hydrogen atoms of the water molecules. Threedimensional differences maps calculated at various stages of the refinement had pronounced maxima, of

height about 0.5 e.Å-3, at the positions expected for the hydrogen atoms of the carbon chain H(1)–H(5) and somewhat smaller maxima representing H(6) and H(7) of the amino group, but only indistinct maxima near the oxygen atoms of the water molecules. Nevertheless, coordinates were assigned to H(8)–H(11) by choosing maxima that were consistent with a reasonable hydrogen-bonding scheme; these coordinates as well as isotropic temperature factors of the hydrogen atoms were included in the final least-squares cycles. The final parameters include a negative temperature factor for H(1) and lead to some rather short N-H and O-H distances; nevertheless, the bond angles and hydrogenbond arrangement are satisfactory and we are confident that the approximate positions of the hydrogen atoms are correctly assigned.

Convergence was first reached at an R index of 0.037. It was noted, however, that the observed values of the structure factors for the stronger reflections were systematically lower than the calculated values. Although this effect is symptomatic of extinction, it could also be the result of systematic errors in scaling – and, in particular, of too low a value for the film factor used to correlate intensities from adjacent films within the same pack. Our empirical value for this factor, averaged over all sets of films and adjusted for the angle of incidence of the diffracted beam, was 3.65 - appreciably lower than the value of about 3.9 usually found for this film (Eastman No-Screen Medical X-ray Film). Accordingly, the data were rescaled, using a value of 3.8 for the film factor, and further least-squares refinement was carried out. Disagreements among the stronger reflections, although still apparent, were much reduced and the refinement was considered complete when no heavy-atom parameter shifted by as much as 0.01 of its standard deviation, no hydrogen coordinate by as much as 0.1 e.s.d., and no hydrogen temperature factor by as much as 0.25 e.s.d.

The observed and calculated structure factors are listed in Table 2, the final heavy-atom parameters and their standard deviations in Table 3, and the hydrogenatom parameters in Table 4. An electron-density projection onto (001), calculated at the conclusion of the refinement, is shown in Fig. 1.

Accuracy of the results

The final R index based on 1023 observed reflections of non-zero weight is 0.032; for none of the reflections too weak to be measured is F_c greater than the threshold value of F_o . The value of

$$\left(\frac{\sum w(F_o^2 - F_c^2)}{m - s}\right)^{\frac{1}{2}}$$

is 0.88. That this latter value is less than 1.0 suggests that the observational errors were over-estimated, perhaps because these errors were derived during the data-reduction process by comparing values of intensities scaled together on the basis of a film factor of 3.65 rather than of the final value 3.80.

Table 2. Observed and calculated structure factors

Within each group are the values of k, $10F_o$, $10F_c$, and α_c (the phase angle). A negative sign preceding $10F_o$ should be read 'less than'. Reflections indicated with an asterisk were assigned zero weight in the final least-squares cycles.

Table 3. The heavy-atom parameters and their standard deviations

Values for the copper atom have been multiplied by 10^5 ; for other atoms, by 10^4 . The temperature factor is in the form $T_i = \exp\{-(b_{11}h^2 + b_{22}k^2 + b_{33}l^2 + b_{12}hk + b_{13}hl + b_{23}kl)_i\}$.

	x	y	z	b_{11}	b_{22}	b_{33}	b_{12}	b_{13}	b_{23}
Cu	77933 (4)	31668 (5)	64970 (7)	315 (4)	455 (5)	786 (10)	3 (7)	-106(10)	-26(11)
C(1)	6595 (3)	5340 (3)	7780 (5)	43 (3)	39 (3)	71 (5)	-15(5)	-15(7)	21 (7)
C(2)	5611 (3)	4661 (3)	6699 (5)	33 (2)	48 (3)	70 (6)	3 (5)	0 (7)	3 (7)
C(3)	4630 (3)	4176 (4)	8015 (5)	31 (3)	58 (4)	95 (7)	3 (5)	10 (6)	20 (8)
C(4)	5067 (3)	3326 (4)	9573 (5)	34 (2)	66(4)	106 (6)	-2(5)	0 (7)	48 (9)
C(5)	4046 (3)	2829 (4)	10756 (5)	35 (2)	50 (3)	97 (6)	-5(5)	12 (7)	-1 (8)
N	6133 (3)	3583 (3)	5602 (4)	38 (2)	57 (3)	78 (5)	-5(4)	-3(6)	-22(7)
O(1)	7666 (2)	4872 (3)	7683 (4)	36 (2)	58 (3)	113 (5)	-9(4)	-20(5)	-35(6)
O(2)	6310 (2)	6286 (2)	8741 (4)	54 (2)	46 (2)	112 (5)	0 (4)	-29(6)	-45(6)
O(3)	4341 (2)	2214 (3)	12249 (4)	33 (2)	64 (3)	98 (5)	5 (4)	13 (5)	49 (6)
O(4)	2969 (2)	3006 (3)	10356 (4)	31 (2)	105 (4)	109 (5)	-5(4)	9 (5)	57 (7)
O(5)	7691 (3)	1344 (3)	5632 (4)	87 (3)	46 (2)	113 (5)	-8(5)	-48 (7)	20 (6)
O(6)	6559 (2)	0959 (3)	2343 (4)	42 (2)	65 (3)	113 (5)	10 (4)	-6(6)	21 (6)

Table 4. Coordinates and standard deviations of the hydrogen atoms

	x	y	z	В
H(1)	0.528(3)	0.522(3)	0.583(5)	-0.7(6)
H(2)	0.419 (5)	0.481 (6)	0.839 (9)	4.1 (14)
H(3)	0.397(5)	0.362 (6)	0.741 (8)	3.9 (14)
H(4)	0.563 (4)	0.393 (5)	1.039 (6)	1.2 (9)
H(5)	0.546(5)	0.241(5)	0.908 (8)	2.5 (12)
H(6)	0.635 (6)	0.361 (8)	0.476 (10)	6.1 (18)
H(7)	0.554(5)	0.282(5)	0.574 (8)	2.5 (12)
H(8)	0.781 (4)	0.084 (5)	0.610 (7)	2.2 (11)
H(9)	0.716 (6)	0.125 (8)	0.431 (10)	5.9 (18)
H(10)	0.699 (6)	0.123 (7)	0.197 (11)	6.2 (20)
H(11)	0.581 (4)	0.129 (4)	0.227 (7)	1·1 (9)

The standard deviations in the coordinates, derived from the residuals and the diagonal elements of the inverse matrix of the final least-squares cycle, correspond to positional uncertainties of about 0.0005 Å for the copper atom, about 0.003–0.004 Å for the C, N and O atoms, and about 0.04–0.08 Å for the hydro-

gen atoms. These uncertainties are unrealistically small for the hydrogen atoms of the amino group and the water molecules, H(6)-H(11), which were poorly represented on the difference maps and show a wide spread in the N-H and O-H distances (Table 5).

The corresponding standard deviations in the bond distances are about 0.004 Å for the Cu-N and Cu-O bonds and about 0.006 Å for the C-C, C-O and C-N bonds. In view of uncertainties in the corrections for thermal motion, which are especially large for the oxygen atoms, we feel that these uncertainties should be increased to, perhaps, 0.008 Å for all the bonds.

Discussion of the results

(i) The molecular structure

The structure as viewed down the c axis is shown in Fig. 2. The coordination about the copper atom is shown in Fig. 3, and the dimensions of the glutamate

Table 5. Bond distances and angles involving the hydrogen atoms

C(2)-H(1)	0·93 Å	C(1)-C(2)-H(1) NC(2)-H(1)	110° 105	C(3)–C(2)–H(1)	110°
C(3)-H(2) C(3)-H(3)	0·86 1·03	C(2)-C(3)-H(2) C(2)-C(3)-H(3) H(2)-C(3)-H(3)	110 115 99	C(4)-C(3)-H(2) C(4)-C(3)-H(3)	113 103
C(4)-H(4) C(4)-H(5)	1·06 1·10	C(3)-C(4)-H(4) C(3)-C(4)-H(5) H(4)-C(4)-H(5)	105 113 117	C(5)-C(4)-H(4) C(5)-C(4)-H(5)	109 101
N-H(6) N-H(7)	0·65 1·03	C(2)-N-H(6) C(2)-N-H(7) H(6)-N-H(7)	127 106 111	Cu-N-H(6) Cu-N-H(7)	89 113
O(5)–H(8) O(5)–H(9)	0·63 1·13	Cu-O(5)-H(8) Cu-O(5)-H(9)	126 112	H(8)-O(5)-H(9)	120
O(6)-H(10) O(6)-H(11)	0·61 0·90	H(10)-O(6)-H(11)	121		

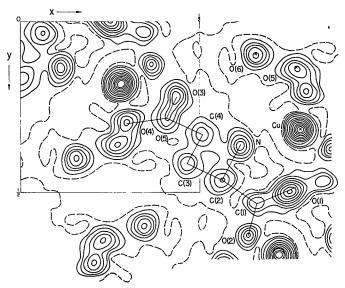


Fig. 1. The final electron-density map projected onto (001). Contours are at 4, 6, $8 \cdots e. A^{-2}$ for the light atoms and 4, 10, 15, $20 \cdots$ for the copper atom: the 2 e.A^{-2} contour is dashed.

group are given in Fig. 4. Distances and angles involving the hydrogen atoms are given in Table 5. Equations of, and deviations from, three planes of interest are given in Table 6.

The coordination about the copper atom is approximately square planar, the square comprising two oxygen atoms and a nitrogen atom of glutamate groups and a water molecule. The average distance from the copper atom to these four ligands is 1.99 Å, slightly longer than the normal single-bond distance of 1.92–1.95 Å (Pauling, 1960). A fifth atom, O(2), occupies an axial position at 2.30 Å from the copper atom; a sixth, O(4), at 2.59 Å and considerably displaced from the other axial position completes a severely distorted octahedron. The four nearest ligands – O(1), O(3), O(5) and N – are coplanar within 0.025 Å; the copper atom is displaced from this plane by 0.15 Å,

in a direction toward the next nearest neighbor O(2) (Table 6).

There are no surprises in the dimensions of the glutamate group (Fig. 4). The lengths of the C(1)-O(1) and C(5)-O(3) bonds approach the values found in carboxylic acids and reflect the fact that these two oxygen atoms are strongly bonded to the copper atom. The two carboxyl groups are planar within 0.01 Å (Table 6). The nitrogen atom is tetra-coordinated and, within experimental error, tetrahedral; it lies only 0.06 Å from the plane of the nearest carboxyl group. The C(2)-C(3)-C(4) and C(3)-C(4)-C(5) angles are significantly larger than tetrahedral, as is usually observed in saturated carbon chains. And as has almost invariably been observed in amino acids and small peptides, the side chain [which begins at C(3)] is in a highly extended configuration.

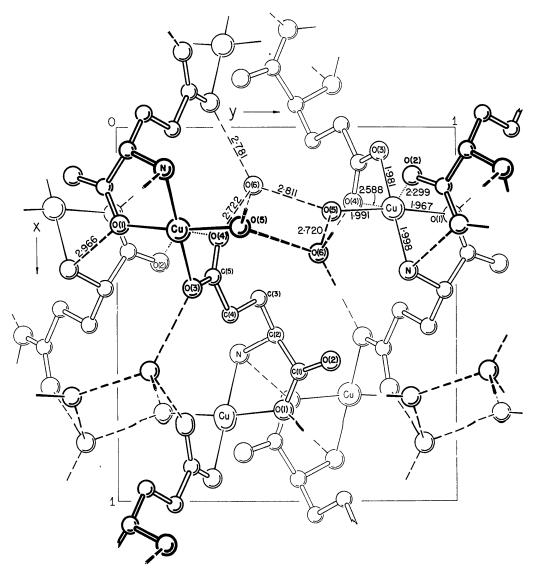


Fig. 2. The structure viewed down the c axis. Hydrogen bonds are indicated by dashed lines and the two long $Cu \cdot \cdot \cdot O$ interactions by dotted lines,

Table 6. Some planes of interest

The coefficients q_i are direction cosines relative to **a**, **b**, and **c**; D is the origin-to-plane distance. Atoms indicated with an asterisk were omitted from the calculation of the least-squares plane; the remaining atoms were given equal weights.

(1) Carboxyl group		
	Atom	Deviation
$q_1 = -0.1858$	C(1)	-0.010
$q_2 = -0.5986$	C(2)	0.003
$q_3 = 0.7791$	O (1)	0.004
$\vec{D} = -0.2686 \text{ Å}$	O(2)	0.004
	*N	-0.056
	*Cu	0.366
(2) Carboxyl group		
$q_1 = 0.0019$	C(4)	0.001
$q_2 = 0.8594$	C(5)	-0.005
$q_3 = 0.5111$	O(3)	0.002
	O(4)	0.002
D = 6.5101 Å	*C(3)	0.179
	*Cu′	0.121
(3) Copper ligands		
$q_1 = -0.4416$	N	-0.023
$q_2 = -0.3403$	O(1)	0.023
$q_3 = 0.8301$	O(5)	0.021
	O(3)'	-0.022
D = -0.8759 Å	*Cu	-0.151
	*O(2)′	-2.439
	*O(4)'	2.067

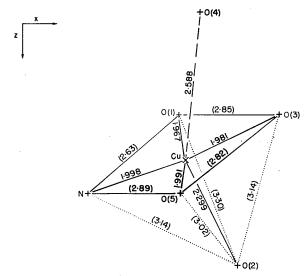


Fig. 3. The coordination about the copper atom.

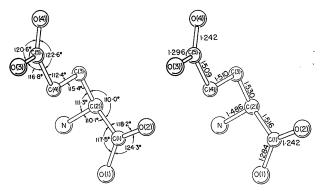


Fig. 4. Bond distances and angles within the glutamate group.

Of the six protons available for hydrogen bonding (two on the nitrogen atom and four on the two water molecules), five participate (Fig. 2). The four $O-H\cdots O$ distances of 2·72 to 2·81 Å and the $N-H\cdots O(1)$ distance of 2·97 are typical values and suggest moderately strong interactions. The second proton on the nitrogen atom, H(7), does not participate in hydrogen bonding, the nearest neighbors in this direction being an O(2) atom at 3·64 Å from the nitrogen atom and an O(3) atom at 3·44 Å.

The lengths and direction cosines of the principal axes of thermal vibration are given in Table 7. The thermal motions are relatively small and not highly anisotropic, reflecting the hardness of the crystals. The directions of maximum motion of O(1) and O(2) make angles of 12° and 29° with the normal to the plane of that carboxyl group, and the directions of maximum motion of C(3), C(4), O(3), and O(4) all lie within 12° from the normal to the second carboxyl group. The

Table 7. Magnitudes and direction cosines, relative to the crystallographic axes, of the principal axes of the thermal ellipsoids

	B_i	$q i^a$	q_{i^b}	q_{i^c}
Cu	$ \left\{ \begin{array}{l} 1.96 \\ 1.77 \\ 1.42 \end{array} \right. $	- 0·090 - 0·596 0·798	-0.980 0.194 0.035	0·175 0·779 0·602
C(1)	$ \begin{cases} 2.44 \\ 1.59 \\ 1.25 \end{cases} $	-0.795 0.606 -0.026	0·485 0·609 - 0·627	0·364 0·511 0·778
C(2)	$ \begin{cases} 2.07 \\ 1.60 \\ 1.46 \end{cases} $	0·155 0·984 0·087	0·985 0·146 - 0·095	0·080 0·102 0·992
C(3)	{ 2.66 1.87 1.47	0·116 0·279 - 0·953	0·895 - 0·445 - 0·021	0·430 0·851 0·301
C(4)	$ \begin{cases} 3.29 \\ 1.75 \\ 1.66 \end{cases} $	-0.022 0.186 -0.982	0·830 -0·545 -0·122	0·558 0·818 0·143
C(5)	$ \begin{cases} 2.21 \\ 2.08 \\ 1.61 \end{cases} $	0·355 0·261 -0·898	-0.843 0.505 -0.187	0·405 0·823 0·399
N	2·56 1·89 1·49	$0.122 \\ -0.966 \\ 0.227$	-0.938 -0.038 0.344	0·324 0·255 0·911
O(1)	$ \begin{cases} 2.95 \\ 2.21 \\ 1.42 \end{cases} $	-0.057 -0.634 0.771	-0.721 0.560 0.407	0·690 0·533 0·489
O(2)	$ \begin{cases} 3.15 \\ 2.42 \\ 1.39 \end{cases} $	-0.625 0.743 0.239	-0.386 -0.560 0.733	0·679 0·366 0·637
O(3)	$ \begin{cases} 3.22 \\ 1.69 \\ 1.46 \end{cases} $	0·128 0·716 - 0·686	0.830 -0.455 -0.321	0·543 0·528 0·653
O(4)		-0.018 0.329 -0.944	0·947 - 0·297 - 0·122	0·320 0·897 0·306
O(5)	$ \begin{cases} 4.57 \\ 2.25 \\ 1.77 \end{cases} $	-0.933 0.337 0.124	0·103 0·581 - 0·807	0·344 0·741 0·577
O(6)	$ \begin{cases} 2.99 \\ 2.32 \\ 1.91 \end{cases} $	0·175 - 0·543 0·822	0.886 - 0.277 - 0.371	0·429 0·793 0·432

principal motion of O(5), a water molecule, is approximately perpendicular to the plane of the three atoms to which it is bonded – Cu, O(6), and O(6'). There is no indication of rigid-body motion of the glutamate group as a whole.

Further discussion of the details of the structure and its relationship to the nearly isostructural zinc compound is given in the following paper (Gramaccioli, 1966).

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The crystal structure of zinc glutamate dihydrate, $ZnC_5H_7NO_4$. $2H_2O$, has been determined and refined by three-dimensional least-squares methods. The crystals are orthorhombic, space group $P2_12_12_1$, with $a=11\cdot190$, $b=10\cdot463$, and $c=7\cdot220$ Å. The final R index for 1035 reflections is 0·032; the standard deviations are about 0·004 Å for the light-atom positions.

Although the compound is nearly isostructural with the corresponding copper compound, the coordination about the zinc atom is a nearly regular square pyramid as compared with a square (with an additional ligand at a considerably greater distance) about the copper atom. The Zn-O and Zn-N distances range from 2.03 to 2.10 Å. This change in coordination has an appreciable effect upon the dimensions of one of the carboxyl groups; otherwise, the dimensions of the glutamate group are not significantly different from those found in copper glutamate. Hydrogen-bond lengths differ by as much as 0.05 Å between the two compounds.

Introduction

The determination of the structure of zinc glutamate (Zn.glu) was originally undertaken as a continuation of studies of the coordination of metal atoms with amino acids and peptides. Interest was further whetted when preliminary photographs showed that crystals of zinc glutamate dihydrate have the same space group, and very nearly the same cell dimensions, as the copper compound Cu.glu (Gramaccioli & Marsh, 1966; hereafter, GM) but that there are differences in intensity of corresponding reflections far greater than would be expected if the two compounds were truly isostructural. Very early in the analysis it became apparent that these differences in intensity were due primarily to differences in the coordinates of the metal atoms, rather than of the amino-acid or water molecules, and that as a result the coordination about zinc would be appreciably dif-

Experimental

White prismatic crystals of zinc glutamate dihydrate, $ZnC_5H_7NO_4.2H_2O$, form on slow evaporation of an aqueous solution of zinc oxide in glutamic acid; like the copper compound, they are elongated along **c** with principal faces {110} and {011}. Unit-cell dimensions were obtained as described by GM, based on measurements of 70 0kl and hk0 reflections from two crystals; the results are summarized in Table 1. Once again, the difference between b values measured for the two crystals is much greater than the apparent experimental uncertainty, although the discrepancy is much less than that observed in Cu.glu. In view of this discrepancy, the e.s.d.'s of the final ('best') values include a lack-of-confidence factor of about 15.

The experimental density was not measured; that calculated on the basis of four molecules per unit cell is 1.937 g.cm⁻³. Absence of reflections h00, 0k0, and 00l with h, k or l odd indicates the space group $P2_12_12_1$.

ferent from that about copper. In order to assure a good comparison between the two compounds, the same care has been taken with the present refinement as with that of the copper compound.

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